# 多成分凝縮系における渦構造 BECからQCDまで



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Topological Quantum Phenomena in Condensed Matter with Broken Symmetries



Keio University 1858 CALAMVS GLADIO FORTIOR

### **References**

### **BEC** (Bose-Einstein condensates)

- Lattice of vortex molecules with Mattia Cipriani [1] Phys.Rev.Lett. 111 (2013) 170401[arXiv:1303.2592 [cond-mat.quant-gas]] [2] Phys.Rev.A88 (2013) 013634 [arXiv:1304.4375 [cond-mat.quant-gas]] Vortex graphs (or N-omers) with M.Eto(衛藤稔)
- [3] Europhys.Lett. 103 (2013) 60006 [arXiv:1303.6048 [cond-mat.quant-gas]]

**QCD** (Quantum Chromodynamics)

- [4] Invited review: Vortices and solitons in dense QCD, with M.Eto, Y.Hirono(広野雄士), S.Yasui(安井繁宏) Prog.Theor.Exp.Phys.:012D01,2014 [arXiv:1308.1535 [hep-ph]]
- [5] Lattice of non-Abelian vortices
  - with M.Kobayashi(小林未知数),E.Nakano(仲野英司) arXiv:1311.2399 [hep-ph]

### **``Pure'' BEC** (99% is condensed)



Cold atomic gases 1995 cold atomic bose gas <sup>87</sup>Rb, <sup>23</sup>Na, <sup>7</sup>Li Cornell (Colorado), Ketterle(MIT) & Wieman (Colorado) 2003 cold atomic fermion gas JILA(Colorado), MIT



doppler laser cooling magneto-optical trap evaporative cooling

Temperature ~  $10^{-6}$ ,  $10^{-7}$  K Number ~  $10^{6}$ , Size ~  $10^{-3}$  cm



### Scalar BEC, <sup>4</sup>He superfluid

Gross-Pitaevskii (nonlinear Schrödinger) Equation

$$i\hbar\frac{\partial\psi}{\partial t} = \left[\frac{\hbar^2}{2M}\nabla^2 + V_{\text{ext}} - \mu + g|\psi|^2\right]\psi = \frac{\delta E}{\delta\psi^*} \qquad g = \frac{4\pi\hbar^2 a_s}{M}$$

Note: derived by Bogoliubov theory weakly interactive Bose gas with point interaction  $V(r) = g\delta(r)$ 



### **Rotation**

Rotating frame 
$$\nabla \rightarrow \nabla - i \frac{M}{\hbar} \Omega \times \mathbf{r}$$
  
 $E[\psi] = \int d^3 \mathbf{r} \left\{ \frac{\hbar^2}{2M} \right| \left( \nabla - i \frac{M}{\hbar} \Omega \times \mathbf{r} \right) \psi \right|^2 + (V_{\text{eff}} - \mu) |\psi|^2 + \frac{g}{2} |\psi|^4 \right\}$   
 $V_{\text{eff}} = V_{\text{ext}} - \frac{M}{2} \Omega^2 r^2$ 

## Superconductors under magnetic fields

= generaction of a vortex lattice

### Vortex lattice in BEC (expretiment), 2001



### MIT [Abo-Shaer et.al, Science 292 (2001) 476]



K. W. Madison et al. PRL 86, 4443 (2001)



**Miscible 2 component BEC** 

$$(\psi_{1},\psi_{2})$$

Gross-Pitaevskii energy functional

$$E[\psi] = \int d^{3}\mathbf{r} \left\{ \sum_{i} \left| \frac{\hbar^{2}}{2m_{i}} |\nabla \psi_{i}|^{2} + (V_{\text{ext}} - \mu_{i}) |\psi_{i}|^{2} \right\} + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$

### **Experiments**

**Rb**, **2** comp **BEC** (of the same atom with different hyperfine  $m_1 = m_2$ )

- (1) |1,-1>, |2, 1>: Matthews, Anderson, Haljan, Hall, Wieman, and Cornell, Phys. Rev. Lett., 83, 2498 (1999).
- (2) (2, 1), (2, 2) : Maddaloni, Modugno, Fort, Minardi and Inguscio, Phys. Rev. Lett. 85, 2413 (2000)









### Vortex lattices of miscible **3** component BECs *without* Rabi $|\Psi_1|^2$ $|\Psi_2|^2$ $|\Psi_{3}|^{2}$ **Cipriani & MN** 10 10 10 Phys.Rev.A88 5 5 5 (2013) 013634 0 0 0 arXiv:1304.4375 -5-5-5[cond-mat.quant-gas] -10-10 -10-10-5 0 5 10 -10 - 5 0 5 10-10 - 5 05 10



Always Abrikosov

# comp = # edges of triangle

Simulating QCD (color superconductor)

### Miscible 2 component BEC + Rabi

Gross-Pitaevskii energy functional

$$E[\psi] = \int d^{3}\mathbf{r} \begin{cases} \sum_{i} \left( \frac{\hbar^{2}}{2m} |\nabla \psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \\ -\sum_{i,j} \omega_{ij} \psi_{i}^{*} \psi_{j} + c.c & \text{internal coherent coupling} \\ (\text{Rabi oscillation}) & \text{Josephson coupling=supercond} \\ -\psi_{1}^{*} \psi_{2} + c.c = -2 |\psi_{1}| |\psi_{2}| \cos(\theta_{1} - \theta_{2}) \end{cases}$$

$$\omega_{12} > 0$$
  $\theta_1 = \theta_2$  Phases coincide

$$\omega_{12} < 0$$
  $\theta_1 = \theta_2 + \pi$   $\pi$  Phase



 $\omega_{12} \neq 0$  $(m_1 = m_2, v_1 = v_2)$  $(\Psi_1, \Psi_2) \sim (e^{i\theta_1}, 1) = e^{i\theta_1/2} (1, 1) e^{i\theta_1\sigma_3/2}$  $(\Psi_1, \Psi_2) \sim (1, e^{i\theta_2}) = e^{i\theta_2/2} (1, 1) e^{-i\theta_2\sigma_3/2}$ b. **b**<sub>2</sub>  $U(1)_{\text{gauge}} U(1)_{\text{relative}}$  $(1,0) = \frac{1}{2}(1,1) + \frac{1}{2}(1,-1): b_1 + r$  $(0,1) = \frac{1}{2}(1,1) - \frac{1}{2}(1,-1): b_2 - r$ Sine-Gordon kink **Kasamatsu**  $- \omega (\psi_1^* \psi_2 + c.c)$ Tsubota & Ueda,  $= -2 |\psi_1| |\psi_2| \omega \cos(\theta_1 - \theta_2)$ PRL93('04) 2 Along the path *r* Kink: Tanaka<sup>(+</sup>01), × Molecule: Babaev('02), Goryo et.al ('07)  $\Delta \theta = \theta_1 - \theta_2$  changes  $2\pi$ for 2 gap superconductors

### Miscible 2 component BEC + Rabi

Gross-Pitaevskii energy functional

$$E[\psi] = \int d^{3}\mathbf{r} \left\{ \sum_{i} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \frac{g_{ij}}{2} |\psi_{i}|^{2} |\psi_{j}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right\} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\nabla\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right) + \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right\} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right\} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} \right\} \right\}$$
  
$$= \int d^{3}\mathbf{r} \left\{ \sum_{i,j} \left( \frac{\hbar^{2}}{2m_{i}} |\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^{2} + (V_{ext} - \mu_{i})|\psi_{i}|^$$

We rotate the system.  $\nabla \rightarrow \nabla - i \Omega \times \mathbf{r}$ We introduce the trapping potential.  $V_{\text{ext}}$ 

















### **Square lattice**



### **Square lattice**











Ichie-Suganuma et.al ('03)









**1.1.1** 

.3



**!**•



Left =absolute minimum Right =local min.

- 0.25 - 0.20

- 0.15

0.10



One can manipulate the shape of graphs as one likes, by changing  $\mathcal{O}_{ij}$ .

7 component BEC

### Eto & MN, Europhys.Lett. 103 (2013) 60006 [arXiv:1303.6048 cond-mat.quant-gas]



### **Quantum Chromo Dynamics (QCD)**

### quarks

Color-flavor locked (CFL) phase

**Quark-Gluon Plasma** 

Liquid-Ga

Nuclear Superfluid

Ь

Temperature

sQGP

Critica

**Hadronic Phase** 

**Quark matter** 



from Fukushima & Hatsuda Rept.Prog.Phys. 74 (2011) 014001

### **Quantum Chromo Dynamics (QCD)**



**Color superconductor** 

$$\Phi_{\alpha i} = \varepsilon_{\alpha\beta\gamma}\varepsilon_{ij\,k}q_{j}^{\beta}q_{k}^{\gamma}$$



$$\alpha = 1,2,3 \ (r,g,b) \ i = 1,2,3 \ (u,d,s)$$

$$\Phi_{\alpha i} = \begin{pmatrix} d_{[g}s_{b]} & s_{[g}u_{b]} & u_{[g}d_{b]} \\ d_{[b}s_{r]} & s_{[b}u_{r]} & u_{[b}d_{r]} \\ d_{[r}s_{g]} & s_{[r}u_{g]} & u_{[r}d_{g]} \end{pmatrix} \overrightarrow{b} = rg$$

$$\overline{u} = ds \quad \overline{d} = sb \quad \overline{s} = ud$$

$$G = U(1)_{B} \times SU(3)_{C} \times SU(3)_{F} \quad \Phi_{\alpha i} \rightarrow e^{i\alpha}g_{color}\Phi_{\alpha i}g_{flavor}$$





lida & Baym, Forbes & Zhitnitsky('02)



Balachandran, Digal & Matsuura (BDM) ('05) Nakano, MN & Matsuura ('07), Eto & MN ('09)















# **Colorful** vortex lattice



$$\Phi_{\alpha i} = \begin{pmatrix} \Delta_{1}(r)e^{i\theta} & 0 & 0 \\ 0 & \Delta_{0}(r) & 0 \\ 0 & 0 & \Delta_{0}(r) \end{pmatrix} \overset{H = SU(3)_{C+F}}{K = [SU(2) \times U(1)]_{C+F}} \\ & & \downarrow \\ K = [SU(2) \times U(1)]_{C+F} \\ & @ \text{ vortex core} \\ \text{Nambu-Goldstone modes localized around the vortex} \\ & & Color \\ C \times \frac{H}{K} = C \times \frac{SU(3)_{C+F}}{SU(2) \times U(1)} \overset{\text{Kelvon magnon}}{= C \times CP^{2}} & Continuous family of solutions exists \\ & = C \times \frac{Color}{SU(2) \times U(1)} \overset{\text{Kelvon magnon}}{= C \times CP^{2}} & Continuous family of solutions exists \\ & = \text{Continuous family} \\ & = \text{Gapless modes propagating along the vortex line} \\ & & L_{Kelvon} = X\dot{Y} - Y\dot{X} - T(X'^{2} + Y'^{2}) \text{ Type-II} \quad \phi = (\phi^{1}, \phi^{2}, \phi^{3}) \text{ C}^{3} \\ & & \mathcal{L}_{\mathbb{C}P^{2}} = C \sum_{\alpha=0,3} K_{\alpha} [\partial^{\alpha} \phi^{\dagger} \partial_{\alpha} \phi + (\phi^{\dagger} \partial^{\alpha} \phi)(\phi^{\dagger} \partial_{\alpha} \phi)], \text{ homogenous } \phi^{\dagger} \phi = 1 \\ \end{array}$$

### **Interaction between vortices**

Long-range vortex-interaction by exchanging phonons



 $\boldsymbol{\chi}$ 

### Positions are locked as Abrikosov's triangular lattice cf) lattice oscillation = Tkachenko mode

CP<sup>2</sup> "color" spin on a triangular lattice

cf) CP<sup>1</sup>=S<sup>2</sup>: Heisenberg spin on a triangular lattice

Short-range vortex-interaction Eto, Hirono, Yasui & MN ('13) by exchanging gluons (massive gauge)  $E_{\text{int,gluon}} \propto G(\phi_1, \phi_2) \exp(-m_{\text{g}}R)$  repulsion among colors  $G(\phi_1, \phi_2) \equiv \phi_1^{\dagger} T^a \phi_1 \ \phi_2^{\dagger} T^a \phi_2 \qquad T^a \ \text{SU(3)}$ Short-range vortex-interaction by exchanging adj scalar (gap) Auzzi-Eto-Vinci('07)  $E_{\text{int,adi}} \propto \Theta(\phi_1, \phi_2) \exp(-m_{\text{adi}}R)$  attraction among colors Kobayashi,Nakano & MN, arXiv:1311.2399  $H = \int dz \sum_{\langle i,j\rangle,A} \left[ -J_{xy} S_{i,A} S_{j,A} \right]$  $E_{\text{int,total}} = E_{\text{int,gluon}} + E_{\text{int,adj}}$ color color ferro anti-ferro  $S_{i,A} := \phi_i^{\dagger} T_A \phi_i, \quad J_{xy} := \Delta^2 G(L)$ 







### Summary

**Vortex Lattices** under rotation

**BEC**: Lattices of vortex molecules

triangular / square lattices,



Rabi (Josephson) interaction, reconnection

**QCD**: Lattices of non-Abelian vortices

superfluid vortex: Abrikosov's triangular lattice

color flux tube

(i) type-I color super  $m_{adj} < m_g$ color ferro = not so colorful lattic

(ii) type-II color super  $m_{\rm g} < m_{\rm adj}$ color anti-ferro = colorful lattice

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